

# Torsion potentials and electronic structure of trifluoromethoxy- and trifluoromethylthiobenzene: an ab initio study

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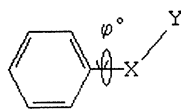
## Abstract

Potential functions of internal rotation about the  $C_{sp^2}-X$  bonds in molecules  $C_6H_5XCF_3$  ( $X = O, S$ ) were calculated at the second-order Møller–Plesset perturbation level of theory with 6-31G(d) basis set. The profile of the potential function and the rotation barrier ( $\Delta E^\ddagger = 3.0$  kJ/mol) found for  $C_6H_5OCF_3$  suggest that, depending on experimental conditions, there can be either free rotation about the  $C_{sp^2}-O$  bond or the conformational equilibrium is shifted to the side of the orthogonal form. The rotational barrier for  $C_6H_5SCF_3$  is 14.7 kJ/mol and the molecule exists in the stable orthogonal conformation. The nature of hybridization, energy and population of lone electron pairs (LPs) on the oxygen and sulfur atoms were considered by using the Natural Bond Orbital (NBO) method. The energy of interactions of the LPs with antibonding  $\pi^*$ -orbitals of the aromatic moiety were estimated for different conformations. The distribution of electron density in the molecules was discussed. The results were compared with analogous calculations on the molecules  $C_6H_5XCH_3$ . © 2002 Elsevier Science B.V. All rights reserved.

**Keywords:** Ab initio; Trifluoromethoxybenzene; Trifluoromethylthiobenzene; Methoxybenzene; Methylthiobenzene; Conformations; NBO analysis

## 1. Introduction

The conformation of trifluoromethoxybenzene and trifluoromethylthiobenzene is characterized by the torsion angle  $\varphi$  between the plane of benzene ring and the  $C_{sp^2}-X-C_{sp^3}$  ( $X = O, S$ ) bond plane:



$\angle\varphi$  from  $0^\circ$  to  $90^\circ$

where  $Y = CH_3$  or  $CF_3$ .

Studies of the molecules by structural physical methods (electron diffraction and X-ray analyses, microwave spectroscopy) have not been made. The first quantum-chemical investigation of the molecule  $C_6H_5OCF_3$  was perfected by the CNDO/2 method [1]. The calculation was carried out with the use of idealized geometry and the torsion angle  $\varphi$  was estimated from the correlation between  $\varphi$  and substituent constants  $\sigma_R^0$ . It was inferred that the molecule exists in the orthogonal conformation due to steric interactions

between the fluorine and *ortho*-hydrogen atoms. Based on analysis of long-range spin–spin coupling constants in NMR spectra of  $C_6H_5OCF_3$  and on ab initio quantum-chemical calculations using minimal STO-3G and split valence 3-21G (with incomplete geometry optimization) basis sets, it was suggested that the molecule has a near-planar structure [2]. The minimum on the potential function of internal rotation about the  $C_{sp^2}-O$  bond is located at  $\varphi = 29.9^\circ$  (STO-3G) or at  $40.1^\circ$  (HF/3-21G). However, according to calculations at the HF/6-31G level, the orthogonal conformation is energetically preferable [2]. It was concluded from the analysis of  $^1H$  and  $^{19}F$  NMR spectra of structurally related compounds *p*- $FC_6H_4OCF_3$  and *p*- $IC_6H_4OCF_3$  in a nematic liquid crystalline solvent that fluorinated anisoles possess the orthogonal conformation [3]. Trifluoromethylthiobenzene was also suggested to have the orthogonal structure on basis of its UV and IR spectra. However, a small amount, about 15%, of the planar conformer was also detected [4]. The X-ray emission  $K_\beta$  spectrum of the sulfur atom in  $C_6H_5SCF_3$  suggests localization of lone electron pairs (LPs) on the sulfur atom [5]. The  $K_\beta$  spectrum is similar to that of *t*- $C_6H_5SC_4H_9$ -*t* [6], which has the orthogonal conformation. Semiempirical quantum-chemical calculations at AM1 [7] and PM3 approximations [8,9] have shown that the  $C_6H_5SCF_3$  molecule has a near-orthogonal conformation. The electronic effect of the  $OCF_3$  and  $SCF_3$  groups on the

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benzene ring is characterized by different substituent constants [10–13], i.e. 0.35 ( $\sigma_m$ ), 0.32 ( $\sigma_p$ ), 0.55 ( $\sigma_I$ ),  $-0.18$  ( $\sigma_R^0$ ) for  $\text{OCF}_3$  and 0.45 ( $\sigma_m$ ), 0.48 ( $\sigma_p$ ), 0.42 ( $\sigma_I$ ), 0.06 ( $\sigma_R^0$ ) for  $\text{SCF}_3$ . These values were used to discuss the influence of the substituents on the reactivity of the molecules [12–17]. From the values of constants  $\sigma_m$  and  $\sigma_p$ , the  $\text{OCF}_3$  and the  $\text{SCF}_3$  groups have large electron-acceptor effect on the aromatic ring. The inductive effect of the groups (see  $\sigma_I$  constants) is also electron-withdrawing and stronger as compared to substituents  $\text{OCH}_3$  ( $\sigma_I = 0.25$  [14]) or  $\text{SCH}_3$  ( $\sigma_I = 0.13$  [15]). In terms of a conjugative effect, the  $\text{OCF}_3$  group is a weak  $\pi$ -donor with respect to the aromatic ring (see  $\sigma_R^0$  constant), much weaker than  $\text{OCH}_3$  ( $\sigma_R^0 = -0.43$  [14]). The  $\pi$ -donor effect of the  $\text{SCF}_3$  group ( $\sigma_R^0 = -0.06$ ) is in contrast to that of  $\text{SCH}_3$  ( $\sigma_R^0 = -0.16$  [15]). Judging by shifts of  $K_\beta$  lines from the sulfur atom in X-ray emission spectra of  $\text{C}_6\text{H}_5\text{SCH}_3$  and  $\text{C}_6\text{H}_5\text{SCF}_3$ , the substitution of hydrogen for fluorine atoms in the methyl group appreciably reduces the electron density on the sulfur atom [18]. The dipole moments of the  $\text{C}_6\text{H}_5\text{OCF}_3$  (2.36 D [13], 2.27 D [19]) and  $\text{C}_6\text{H}_5\text{SCF}_3$  (2.50 D [19]) molecules are greater than for  $\text{C}_6\text{H}_5\text{OCH}_3$  (1.24 D [20]) and  $\text{C}_6\text{H}_5\text{SCH}_3$  (1.31 D [21]). Some physicochemical characteristics and details on synthesis of  $\text{C}_6\text{H}_5\text{XCF}_3$  are presented elsewhere [12,13,22–29].

The aim of the present work is to calculate potential functions of internal rotation about the  $\text{C}_{\text{sp}^2}\text{--O}$  and  $\text{C}_{\text{sp}^2}\text{--S}$  bonds in  $\text{C}_6\text{H}_5\text{OCF}_3$  and the  $\text{C}_6\text{H}_5\text{SCF}_3$ , to determine stationary (minimum and maximum) points, geometric parameters and to clarify the nature of intramolecular interactions. These are given for comparison of the corresponding data obtained for the  $\text{C}_6\text{H}_5\text{OCH}_3$  and  $\text{C}_6\text{H}_5\text{SCH}_3$  molecules at the same level of theory.

## 2. Computational details

The ab initio computations at the second level Møller–Plesset perturbation theory (MP2) [30,31] have been carried out with the use of Gaussian 98W (Revision A.7) program [32]. All orbitals were active. The standard double-split valence basis set 6-31G(d) with Cartesian (6d) representation of the d functions was employed [33]. The standard convergence criteria were used. Torsional potentials for the molecules were constructed by constraining the torsion angle  $\varphi$  to fixed values 0, 15, 30, 45, 60, 75 and 90° and

optimizing all other geometrical parameters. The positions of stationary points on PES were refined by re-optimization of all geometric parameters including the torsion angle  $\varphi$ . Vibrational frequencies were calculated numerically throughout. Stationary points on PES were verified as equilibrium geometries by the absence of imaginary frequencies, and transition states were characterized by the presence of a single imaginary frequency. Natural Bond Orbital (NBO) [34,35] population analysis was performed with the use of the NBO program, Version 3.1 (link 607, Gaussian 98W) [36].

## 3. Conformations

The total energies ( $E_{\text{tot}}$ , a.u.) for the  $\text{C}_6\text{H}_5\text{OCF}_3$  and  $\text{C}_6\text{H}_5\text{SCF}_3$  molecules are displayed in Table 1. The energies of the rotamers ( $\Delta E$ , kJ/mol) relative to the minimum of the potential function of the internal rotation about the  $\text{C}_{\text{sp}^2}\text{--X}$  bonds are given in parentheses. Such a minimum for the  $\text{C}_6\text{H}_5\text{OCF}_3$  molecule is in the domain of the orthogonal conformation, whereas the maximum lies in the region of the planar conformation (see Fig. 1). The rotation barrier, calculated without ZPVE correction with accuracy determined by the 15°-increment, is 4.35 kJ/mol. The Hessian calculations revealed a negative eigenvalue ( $-24.64 \text{ cm}^{-1}$ ) at the point corresponding to the conformation  $\varphi = 0^\circ$ . Therefore, this conformation is the maximum. The conformation with the torsion angle  $\varphi = 90.1^\circ$  ( $E_{\text{tot}} = -642.7953837$  a.u.) is the minimum. The Hessian matrix has only positive eigenvalues at this point. The ZPVE correction is 0.106712 hartree per particle (a scaling factor of is used 0.9661 [37]) at the point of maximum and 0.107241 hartree per particle at the point of minimum. The ZPVE correction decreases the rotation barrier to 3.04 kJ/mol. The energy barrier  $\Delta E^\ddagger < kT$  (at RT = 2.5 kJ/mol) corresponds to free internal rotation [38]. Thus, a “border” situation is realized in the  $\text{C}_6\text{H}_5\text{OCF}_3$  molecule. The molecular moieties have free rotation about the  $\text{C}_{\text{sp}^2}\text{--O}$  bond or restricted motion with larger amplitude in the potential well with a broad flat bottom. The molecule is a conformationally non-rigid system. However, the conformational equilibrium is shifted to the orthogonal form. Note that experimental conditions can affect the conformational equilibrium.

Table 1

The total energies ( $-E_{\text{tot}}$ , a.u.) for  $\text{C}_6\text{H}_5\text{XCH}_3$  and  $\text{C}_6\text{H}_5\text{XCF}_3$  (X = O and S)<sup>a</sup>

Compound	$\varphi^\circ$							
	0	15	30	45	60	75	90	
$\text{C}_6\text{H}_5\text{OCH}_3$	345.6831291 (0.00)	.6827475 (1.00)	.6817110 (3.72)	.6804153 (7.13)	.6796587 (9.11)	.6795371 (9.43)	.6795188 (9.48)	
$\text{C}_6\text{H}_5\text{SCH}_3$	668.3024735 (2.88)	.3025569 (2.66)	.3027167 (2.24)	.3029312 (1.67)	.3033128 (0.16)	.3035414 (0.07)	.3035681 (0.00)	
$\text{C}_6\text{H}_5\text{OCF}_3$	642.7937269 (4.35)	.7938043 (4.15)	.7939728 (3.70)	.7941498 (3.24)	.7944747 (2.39)	.7950598 (0.85)	.7953834 (0.00)	
$\text{C}_6\text{H}_5\text{SCF}_3$	965.3850386 (15.23)	.3857070 (13.60)	.3870953 (9.95)	.3885169 (6.22)	.3897743 (2.92)	.3906761 (0.55)	.3908858 (0.00)	

<sup>a</sup> The energies of the rotamers ( $\Delta E$ , kJ/mol) relative to the minimum are given in parentheses.

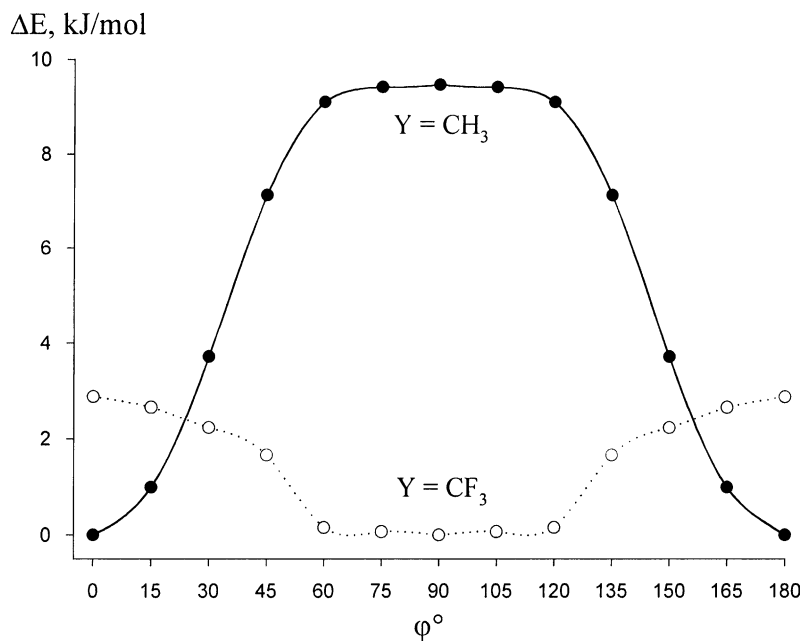


Fig. 1. Potential functions (at MP2/6-31G(d) level) for internal rotation about the  $C_{sp^2}-O$  bond in  $C_6H_5OY$  ( $Y = CH_3$  or  $CF_3$ ).

The energy minimum of the potential function of internal rotation about the  $C_{sp^2}-S$  bond in the  $C_6H_5SCF_3$  molecule also lies close to the orthogonal conformation and the maximum is close to planar (see Fig. 2). The uncorrected barrier is 15.23 kJ/mol. The Hessian for conformation with the torsion  $\varphi = 0^\circ$  has a negative eigenvalue ( $-54.46 \text{ cm}^{-1}$ ), therefore, it is the saddle point. The minimum corresponds to the conformation with  $\varphi = 90^\circ$ . The ZPVE corrections for the points of maximum and minimum are 0.103337 and 0.103550 hartree per particle, respectively. The corrected rotational barrier is 14.67 kJ/

mol. Thus, the  $C_6H_5SCF_3$  molecule exists in the orthogonal conformation.

The planar conformation of the  $C_6H_5OCH_3$  molecule is the energy minimum of the potential function of internal rotation about the  $C_{sp^2}-O$  bond, and the orthogonal form is the maximum (Fig. 1). This is in agreement with the results of electron diffraction [39], microwave spectroscopy [40] and quantum-chemical calculation at MP2/6-31G(d)//HF/6-31G(d) [41,42] and MP2/6-311 + G(d) [43] levels. The rotational barrier of the  $C_6H_5OCH_3$  molecule, without ZPVE correction, is 9.48 kJ/mol. At the point of energy

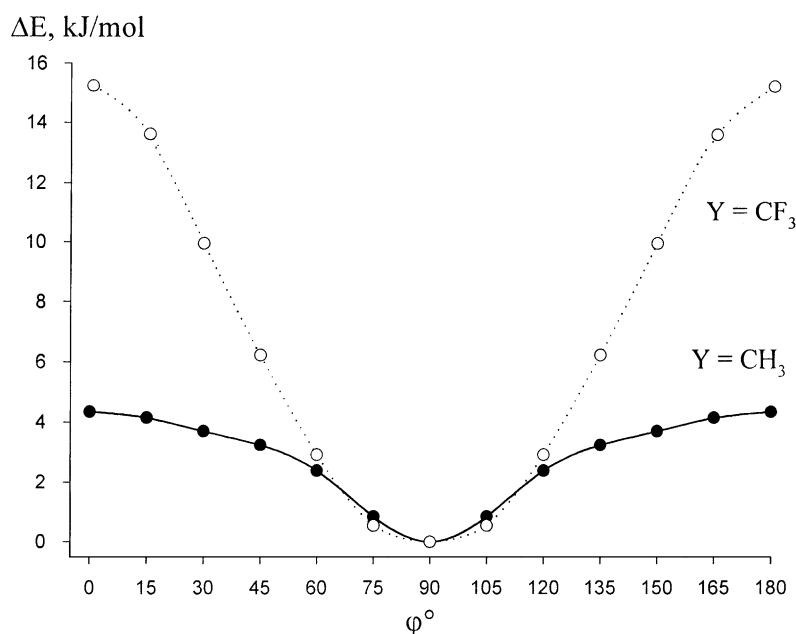


Fig. 2. Potential functions (at MP2/6-31G(d) level) for internal rotation about the  $C_{sp^2}-S$  bond in  $C_6H_5SY$  ( $Y = CH_3$  or  $CF_3$ ).

minimum ( $\varphi = 0^\circ$ ), the Hessian has all positive eigenvalues and at the point of maximum ( $\varphi = 90^\circ$ ) one negative eigenvalue ( $-45.56 \text{ cm}^{-1}$ ). The ZPVE corrections are  $0.129984$  hartree per particle ( $\varphi = 0^\circ$ ) and  $0.129339$  ( $\varphi = 90^\circ$ ). The corrected rotational barrier is  $7.78 \text{ kJ/mol}$ , which is in a good agreement with the value of  $7.91 \text{ kJ/mol}$  found at MP2/6-311 + G(d, p) [43] level of theory.

The energy minimum on the potential function of internal rotation about the  $\text{C}_{\text{sp}^2}\text{-S}$  bond in the  $\text{C}_6\text{H}_5\text{SCH}_3$  molecule is at the near-orthogonal conformation region and the planar conformation is the maximum (Fig. 2). The uncorrected rotation barrier is  $2.88 \text{ kJ/mol}$ . The complete geometrical optimization shows that the energy minimum is a conformer with the torsion  $\varphi = 85.6^\circ$  ( $E_{\text{tot}} = -668.3035682 \text{ a.u.}$ ). At this point, the Hessian matrix has only positive eigenvalues. The ZPVE correction for the total energy is  $0.126353$  hartree per particle. The major maximum is observed at the torsion angle  $\varphi = 0^\circ$ . The Hessian matrix at this point has one negative eigenvalue ( $-29.86 \text{ cm}^{-1}$ ). The correction is  $0.126433$  hartree per particle. The ZPVE corrected barrier between the planar and non-planar (near-orthogonal) forms is  $3.08 \text{ kJ/mol}$ . The Hessian matrix for the structure with the torsion angle  $\varphi = 90^\circ$  has one negative eigenvalue ( $-30.05 \text{ cm}^{-1}$ ) and therefore can be identified as the point of the second maximum. The barriers between these two non-planar forms at  $\varphi = 90^\circ$  is assessed at the purely formal value of  $0.02 \text{ kJ/mol}$  which is comparable with the calculation error. Thus, the  $\text{C}_6\text{H}_5\text{SCH}_3$  molecule, like  $\text{C}_6\text{H}_5\text{OCF}_3$ , is a conformationally non-rigid molecule. Experimental conditions can substantially influence the conformational equilibrium and the profile of the potential function. Physical methods on non-rigid molecules usually reveal only a so-called “effective” conformation, i.e. a hypothetical three-dimensional structure averaged over all possible conformations, having regard to their statistical weights. Electron diffraction identified the conformation of  $\text{C}_6\text{H}_5\text{SCH}_3$  in the vapour phase having a torsion angle  $\varphi = 45 \pm 10^\circ$  [44]. When free rotation about the  $\text{C}_{\text{sp}^2}\text{-S}$  bond occurs, such a value of  $\varphi$  reflects the “effective” conformation that is the average of  $0$  and  $90^\circ$ . This inference is confirmed by MP4(SDQ)/6-31G(d, p)//HF/6-31G(d, p) calculations on the  $\text{C}_6\text{H}_5\text{SCH}_3$  molecule which estimate the rotational barrier at  $2.19 \text{ kJ/mol}$  [45]. At the same time, the values of *Kerr* constants suggest that in solution ( $\text{CCl}_4$ ,  $\text{C}_6\text{H}_6$ ) the conformational equilibrium is shifted to the side of a flatter structure ( $\varphi = 23 \pm 5^\circ$ ) [21].

Thus, the orthogonal or near-orthogonal conformers are the energy minima on the potential function of internal rotation about the  $\text{C}_{\text{sp}^2}\text{-X}$  bonds for  $\text{C}_6\text{H}_5\text{OCF}_3$ ,  $\text{C}_6\text{H}_5\text{SCF}_3$  and  $\text{C}_6\text{H}_5\text{SCH}_3$  and the major maxima are planar forms.  $\text{C}_6\text{H}_5\text{OCH}_3$  has a planar structure.  $\text{C}_6\text{H}_5\text{OCH}_3$  and  $\text{C}_6\text{H}_5\text{SCF}_3$  exist in stable conformations while the  $\text{C}_6\text{H}_5\text{OCF}_3$  and  $\text{C}_6\text{H}_5\text{SCH}_3$  molecules belong to conformationally non-rigid structures. The conformation of the  $\text{C}_6\text{H}_5\text{XCF}_3$  and  $\text{C}_6\text{H}_5\text{XCH}_3$  molecules is determined by two major factors, by the effect of  $n,\pi$ -conjugation of the

LP of the O and S atoms with the  $\pi$ -system of the aromatic ring, which increases stability of the planar forms, and by steric interactions between the  $\text{CF}_3$  and  $\text{CH}_3$  groups and the *ortho*-H atoms of the aromatic ring, which decreases stability of less strained non-planar forms. The  $\text{CF}_3$  group is larger than  $\text{CH}_3$  as follows, for instance, from the comparison of Van der Waals radii of hydrogen ( $1.20 \text{ \AA}$ ) and fluorine ( $1.47 \text{ \AA}$ ) atoms [46,47] or steric constants  $E_s^\circ$  of these groups, i.e.  $-1.24$  ( $\text{CH}_3$ ) and  $-2.40$  ( $\text{CF}_3$ ) [47].

Therefore, with the same X atoms, the planar conformations of the  $\text{F}_3\text{CX}$ -substituted benzenes are sterically more crowded than for the  $\text{H}_3\text{CX}$ -substituted molecules. In the series of the compounds with different atoms X (O or S) but identical groups  $\text{CH}_3$  or  $\text{CF}_3$ , the molecules with X = O have more flattened conformations and lower barriers to rotation about the  $\text{C}_{\text{sp}^2}\text{-X}$  bonds owing to more effective  $n,\pi$ -conjugation of the oxygen LP with the aromatic ring compared to sulfur atom.

#### 4. Geometrical parameters

Table 2 reports the optimized values of valence angles and bond lengths for the  $\text{C}_6\text{H}_5\text{OCF}_3$  and  $\text{C}_6\text{H}_5\text{SCF}_3$  molecules computed as the torsion angle varied from  $0$  to  $90^\circ$ .

It is likely that the CXC and CCX valence angles should depend noticeably on conformation. The lengths of the  $\text{C}_{\text{sp}^2}\text{-X}$  and  $\text{C}_{\text{sp}^3}\text{-X}$  bonds vary with conformation to a lesser extent. Interestingly, the  $\text{C}_{\text{sp}^2}\text{-O}$  bond in  $\text{C}_6\text{H}_5\text{OCF}_3$  is longer than the  $\text{C}_{\text{sp}^3}\text{-O}$  bond. The reason for this might be the high polarity of the triad  $\text{C}_i\text{-O-C}_F$  that is reflected in substantial difference of charges on the atoms. The  $\text{C}_6\text{H}_5\text{SCF}_3$  molecule shows the usual trend: the  $\text{C}_{\text{sp}^2}\text{-S}$  bond is shorter than the  $\text{C}_{\text{sp}^3}\text{-S}$ . The C–F bond in  $\text{C}_6\text{H}_5\text{OCF}_3$  is much shorter than in  $\text{C}_6\text{H}_5\text{SCF}_3$  (see values of  $\sum l(\text{C-F})/3$ ) in Table 2. The length of the C–X bond in  $\text{C}_6\text{H}_5\text{XCH}_3$  varies depending on conformation, in the ranges:  $1.352\text{--}1.355 \text{ \AA}$  ( $\text{C}_{\text{sp}^2}\text{-O}$ ),  $1.422\text{--}1.429 \text{ \AA}$  ( $\text{C}_{\text{sp}^3}\text{-O}$ ),  $1.770\text{--}1.781 \text{ \AA}$  ( $\text{C}_{\text{sp}^2}\text{-S}$ ) and  $1.802\text{--}1.813 \text{ \AA}$  ( $\text{C}_{\text{sp}^3}\text{-S}$ ). Thus, the bonds ( $\text{C}_{\text{sp}^2}\text{-X}$ ) in  $\text{C}_6\text{H}_5\text{XCF}_3$  are longer and the bonds ( $\text{C}_{\text{sp}^3}\text{-X}$ ) are shorter than in the  $\text{C}_6\text{H}_5\text{XCH}_3$  molecules with the same X atoms. The difference is particularly appreciable for compounds with X = O.

The valence angles in the  $\text{C}_6\text{H}_5\text{XCH}_3$  molecules change in the ranges  $117\text{--}112^\circ$  ( $\angle\text{COC}$ ),  $125\text{--}120^\circ$  ( $\angle\text{CCO}$ ),  $103\text{--}99^\circ$  ( $\angle\text{CSC}$ ) and  $125\text{--}120^\circ$  ( $\angle\text{CCS}$ ). The CSC, CCO and CCS angles in all the molecules considered are rather similar and the COC angle in  $\text{C}_6\text{H}_5\text{OCF}_3$  is slightly greater than in  $\text{C}_6\text{H}_5\text{OCH}_3$  at corresponding torsion angles  $\varphi$ .

#### 5. Intramolecular interactions

To analyse the intramolecular interactions, the wavefunctions obtained for the  $\text{C}_6\text{H}_5\text{XCF}_3$  and  $\text{C}_6\text{H}_5\text{XCH}_3$  molecules were represented as orbitals corresponding to traditional

Table 2  
Variation of bond angles ( $^{\circ}$ ) and lengths ( $\text{\AA}$ ) in  $\text{C}_6\text{H}_5\text{OCF}_3$  and  $\text{C}_6\text{H}_5\text{SCF}_3$

Compound	$\varphi^{\circ}$						
	0	15	30	45	60	75	90
$\text{C}_6\text{H}_5\text{OCF}_3$							
$\angle\text{COC}$	119.93	119.71	119.06	118.11	116.94	115.56	114.77
$\angle\text{CCO}$	125.08	124.80	124.04	124.02	121.86	120.36	118.93
$l(\text{C}_{\text{sp}^2}-\text{O})$	1.402	1.402	1.403	1.405	1.408	1.409	1.409
$l(\text{C}_{\text{sp}^3}-\text{O})$	1.352	1.352	1.353	1.354	1.354	1.354	1.355
$\sum l(\text{C}-\text{F})/3$	1.344	1.344	1.343	1.344	1.344	1.343	1.343
$\text{C}_6\text{H}_5\text{SCF}_3$							
$\angle\text{CSC}$	103.02	102.57	101.57	110.35	99.10	97.70	97.01
$\angle\text{CCS}$	125.97	125.48	124.32	123.01	121.72	120.43	119.80
$l(\text{C}_{\text{sp}^2}-\text{S})$	1.781	1.781	1.780	1.781	1.781	1.782	1.782
$l(\text{C}_{\text{sp}^3}-\text{S})$	1.794	1.794	1.796	1.798	1.798	1.798	1.797
$\sum l(\text{C}-\text{F})/3$	1.347	1.347	1.347	1.347	1.347	1.347	1.347

chemical terminology within the framework of the NBO technique. The energy and the hybridization state of the heteroatom LPs were determined. LPs of the O and S atoms are non-equivalent. One of the pairs (LP1) comprises a hybridized orbital, as in accord with the NBO analysis the s-character of the LP1 is assessed at about 40% for the oxygen and 70% for the sulfur atom. With the same X atoms, the contribution of s-AO into LP1 depends only little on the nature of the  $\text{CH}_3$  or  $\text{CF}_3$  groups. The second lone electron pair (LP2) of the O and S atoms is practically pure p-AO. As a result, the LP1 of the O and S atoms noticeably differ in energy. Shown in Fig. 3 are the values of energy of LPs ( $E_{\text{LP}}$  in eV) in the conformationally non-rigid molecules ( $\text{C}_6\text{H}_5\text{OCF}_3$  and  $\text{C}_6\text{H}_5\text{SCH}_3$ ) throughout the range of the torsion angle  $\varphi$  and in the conformationally stable molecules ( $\text{C}_6\text{H}_5\text{SCF}_3$  and  $\text{C}_6\text{H}_5\text{OCH}_3$ ) only at the points of minimum of the potential function. For compounds with the  $\text{CF}_3$  group  $E_{\text{LP1}}$  has greater magnitude than for the  $\text{CH}_3\text{X}$ -substituted molecules. With the same groups  $\text{CF}_3$  or  $\text{CH}_3$ , the  $E_{\text{LP1}}$  is little dependent on the nature of the X atom. The influence of X atoms and  $\text{CF}_3$  or  $\text{CH}_3$ , on the value of the  $E_{\text{LP2}}$  is more substantial. For compounds with the  $\text{OCF}_3$  or  $\text{OCH}_3$  groups, the values of  $E_{\text{LP2}}$  are located lower (in energy scale) than for their sulfur analogs. With the same X, for the compounds with the  $\text{CF}_3$  group the energy of  $E_{\text{LP2}}$  is lower. Thus, the nature of the  $\text{CF}_3$  or  $\text{CH}_3$  groups has the major influence on the difference in energy of the LP1's, whereas the substituents interact mainly with the  $\sigma^*$ -type orbitals. The LP2's interact chiefly with  $\pi^*$ -orbitals and the differences in the energies of LP2's are mainly determined by the nature of the heteroatom X.

The interaction of the lone pairs of O and S with the aromatic ring and the  $\text{CF}_3$  and  $\text{CH}_3$  groups affect their population (see values of  $E_{\text{LP}}$  in Fig. 3). When passing from the planar to the orthogonal conformations, the population of LP1 is decreasing, while that of LP2 is increasing. The conformation-dependent changes in the population of LP2, are due to the steric disruption of the  $n,\pi$ -conjugation

between the LP and the ring  $\pi$ -system. That is, with lowering of coplanarity, the delocalization of LP2 onto the aromatic moiety decreases and its localization on the heteroatom increases. However, the population of LP1 and LP2 cannot be a measure of the  $n,\pi$ -interaction with the ring as the lone pairs also interact with the antibonding  $\sigma^*$ -orbitals of C–F or C–H.

Table 3 lists the values of  $E_{n,\pi}$  of the energy of donor–acceptor interaction of the LP1 and LP2 on the X atoms with antibonding  $\pi^*$ -orbitals of the aromatic ring. The calculations show that the hybridized lone pair (LP1) in planar and near-planar conformations ( $\varphi = 0\text{--}30^{\circ}$ ) does not interact with the ring  $\pi^*$ -orbitals. However, for torsion angles  $\varphi = 30\text{--}45^{\circ}$ , such an interaction starts to occur and peaks in the orthogonal form. In the conformation with  $\varphi = 90^{\circ}$  the values of  $E_{n,\pi}$  (LP1) for anisoles are four to five times greater than for thioanisoles evidently due to higher contribution of the p-component into LP1 of the oxygen as compared with the sulfur atom. LP2 interacts much more strongly with the antibonding  $\pi^*$ -orbitals of the aromatic ring in the planar and near-planar conformations, but in the orthogonal form such an interaction is absent as the p-orbital is displaced from the plane of the aromatic  $\pi$ -system. With the same conformations (excluding the orthogonal one), the value of  $E_{n,\pi}$  (LP2) decreases in the sequence:  $\text{C}_6\text{H}_5\text{OCH}_3 > \text{C}_6\text{H}_5\text{OCF}_3 > \text{C}_6\text{H}_5\text{SCH}_3 > \text{C}_6\text{H}_5\text{SCF}_3$ . That is with the same X atoms and in the same conformation the substitution of the  $\text{CF}_3$  for  $\text{CH}_3$  groups reduces the interaction energy of LP2 with the antibonding  $\pi^*$ -orbitals. The reason may be the competitive donor–acceptor interactions of the LP2 with antibonding  $\sigma^*$ -orbitals of the C–F bonds. As the NBO analysis shows, such an interaction is possible with two C–F bonds pointing toward the benzene ring and it is calculated to be about 67–87 kJ/mol for  $\text{OCF}_3$  and about 40–60 kJ/mol for the  $\text{SCF}_3$  group. A similar interaction of LP2 with  $\sigma^*$ -orbitals of the C–H bonds is observed in the  $\text{H}_3\text{CX}$ -substituted molecules, but its energy is appreciably lower, i.e. 25–30 kJ/mol for the  $\text{OCH}_3$  group

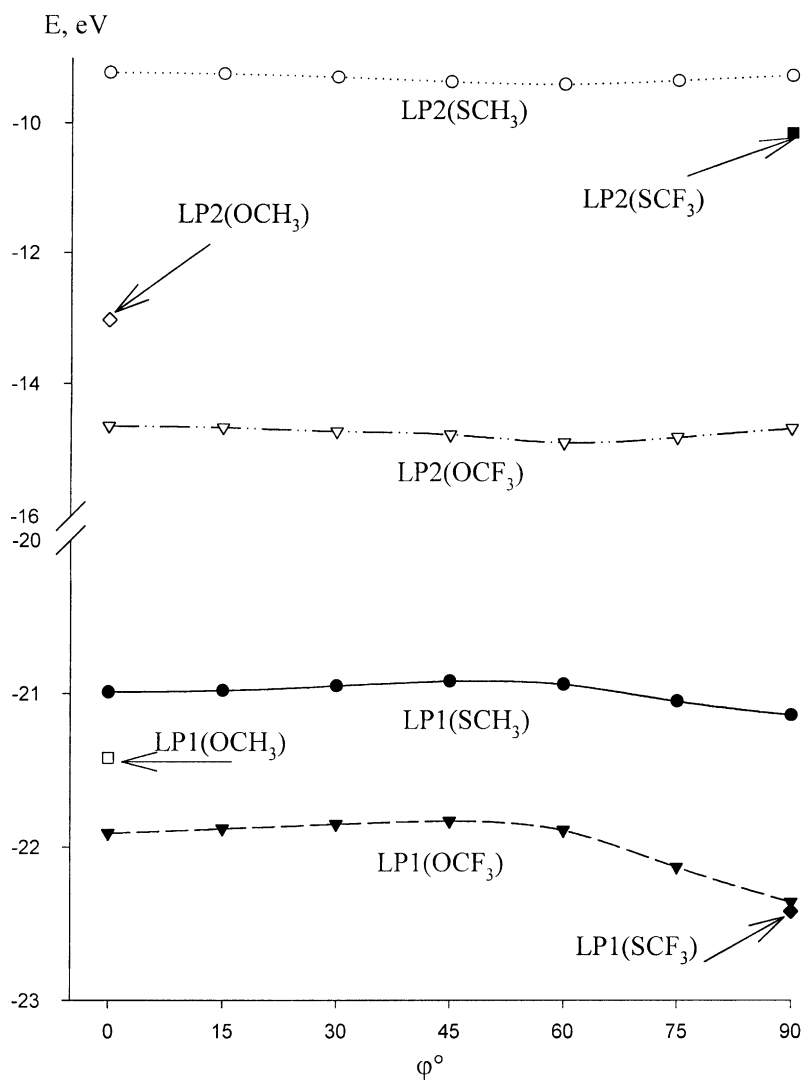
Fig. 3. The energies of LPs ( $E_{LP}$  in eV) at MP2/6-31G(d) level.

Table 3

Population ( $P$ , electron) and energy of LP,  $\pi^*$ -interaction ( $E$ , kJ/mol) for  $C_6H_5XCH_3$  and  $C_6H_5XCF_3$  ( $X = O$  or  $S$ )

XY	Parameters	$\varphi^\circ$							
		0	15	30	45	60	75	90	
OCH <sub>3</sub>	$P(LP1)$	1.94076	1.94050	1.93958	1.93745	1.93332	1.92778	1.92321	
OCF <sub>3</sub>	$P(LP1)$	1.95333	1.95305	1.95219	1.95060	1.94809	1.94554	1.94428	
SCH <sub>3</sub>	$P(LP1)$	1.95307	1.95305	1.95286	1.95223	1.95119	1.95009	1.94952	
SCF <sub>3</sub>	$P(LP1)$	1.97839	1.97843	1.97844	1.97812	1.97738	1.97675	1.97649	
OCH <sub>3</sub>	$P(LP2)$	1.83212	1.83583	1.84672	1.86419	1.88456	1.90081	1.90926	
OCF <sub>3</sub>	$P(LP2)$	1.88726	1.88873	1.89309	1.89978	1.90765	1.91468	1.91795	
SCH <sub>3</sub>	$P(LP2)$	1.84065	1.84460	1.85596	1.87290	1.88963	1.90096	1.90515	
SCF <sub>3</sub>	$P(LP2)$	1.88206	1.88434	1.89050	1.89923	1.90862	1.91578	1.91861	
OCH <sub>3</sub>	$E_{LP1,\pi}$	0.0	0.0	0.0	4.7	8.8	17.3	30.1	
OCF <sub>3</sub>	$E_{LP1,\pi}$	0.0	0.0	0.0	8.0	14.9	23.2	28.3	
SCH <sub>3</sub>	$E_{LP1,\pi}$	0.0	0.0	0.0	2.3	3.9	5.7	6.6	
SCF <sub>3</sub>	$E_{LP1,\pi}$	0.0	0.0	0.0	2.6	4.0	5.2	5.7	
OCH <sub>3</sub>	$E_{LP2,\pi}$	146.9	139.5	118.9	88.4	54.4	25.9	0.0	
OCF <sub>3</sub>	$E_{LP2,\pi}$	105.5	99.9	84.3	61.8	35.9	11.9	0.0	
SCH <sub>3</sub>	$E_{LP2,\pi}$	104.0	96.8	77.4	50.3	26.5	7.8	0.0	
SCF <sub>3</sub>	$E_{LP2,\pi}$	90.4	83.2	60.3	41.8	21.7	6.4	0.0	

and 21–23 kJ/mol for the SCH<sub>3</sub> group. The interaction of LP1 with  $\sigma^*$ -orbitals of the C–F and C–H bonds is 4–8 kJ/mol (OCH<sub>3</sub> SCH<sub>3</sub>, SCF<sub>3</sub>) and 34–35 kJ/mol (OCF<sub>3</sub>) and these energies are little dependent on the structure of the molecules.

Consequently, charge is transferred from the oxygen atom onto the aromatic ring in planar C<sub>6</sub>H<sub>5</sub>OCH<sub>3</sub> owing to interaction of the LP2 with antibonding  $\pi^*$ -orbitals of the ring ( $E_{n,\pi} \approx 147$  kJ/mol). For C<sub>6</sub>H<sub>5</sub>OCH<sub>3</sub> with torsion angle  $\varphi$  varies from 45° (free rotation) to 90° and for this reason the donor effect of the LP2 on the ring is substantially weakened. If the conformational equilibrium is shifted to the orthogonal form, the LP2 does not interact with the antibonding  $\pi^*$ -orbitals of the ring, but at the same time donor interaction occurs between the LP1 and the  $\pi^*$ -orbitals ( $E_{n,\pi} \approx 28$  kJ/mol). C<sub>6</sub>H<sub>5</sub>SCF<sub>3</sub> is orthogonal, where the interaction of LP1 with the antibonding  $\pi^*$ -orbitals is weak ( $E_{n,\pi} \approx 6$  kJ/mol). In the conformation of C<sub>6</sub>H<sub>5</sub>SCH<sub>3</sub> with the torsion angle  $\varphi = 45^\circ$ , the overall donor interaction of LP1 and LP2 with antibonding  $\pi^*$ -orbitals of the ring is estimated at about 53 kJ/mol. In the orthogonal conformation, only LP1 interact with the  $\pi^*$ -orbitals ( $E_{n,\pi} \approx 7$  kJ/mol).

The fluorine atom has three lone electron pairs of which one (LP<sub>F1</sub>) is represented according to the NBO analysis, by the hybrid sp-orbital (about 70% s-character) and two others (LP<sub>F2</sub> and LP<sub>F3</sub>) are considered as practically pure p<sub>y</sub>- and p<sub>z</sub>-AO. In each fluorine atom LP2 interacts with the antibonding  $\sigma^*$ -orbitals of the X–C<sub>F</sub> bond with the energy 57–59 kJ/mol in the C<sub>6</sub>H<sub>5</sub>OCF<sub>3</sub> molecule ( $\varphi = 45\text{--}90^\circ$ ) and of 37–43 kJ/mol in C<sub>6</sub>H<sub>5</sub>SCF<sub>3</sub> ( $\varphi = 90^\circ$ ). The lone pairs of the fluorine interact also with antibonding  $\sigma^*$ -orbitals of the neighbouring C–F bonds. In the CF<sub>3</sub> group, two C–F bonds (*syn*) point toward the benzene ring and the third C–F bond (*anti*) points in the opposite direction. In the planar and the orthogonal conformation, both *syn*-fluorine atoms are

equivalent. At torsion angles  $0^\circ < \varphi < 90^\circ$ , the CF<sub>3</sub> group turns and the *syn*-atoms become non-equivalent. Depending on the spatial orientation of antibonding  $\sigma^*$ -orbitals of the C–F bonds with respect to the lone pairs on the neighbouring fluorine atoms, the energy of their interaction changes substantially and in particular situations exceeds 120 kJ/mol.

## 6. Charge distribution

The use of the set of natural aromatic orbitals (NAO) constructed in an arbitrary basis of AO allows compensation for the disadvantages of the Mulliken population analysis [48]. Table 4 lists the natural charges ( $q$ ) on atoms of the C<sub>6</sub>H<sub>5</sub>XCF<sub>3</sub> and C<sub>6</sub>H<sub>5</sub>XCH<sub>3</sub> molecules in conformations with torsion angles of  $\varphi = 0, 45$  and  $90^\circ$ .

It is convenient to consider the influence of the XCF<sub>3</sub> and XCH<sub>3</sub> groups on electron density distribution in the aromatic ring by comparing the differences of charges on the atoms in the molecules and in unsubstituted benzene, i.e.  $\Delta q = q(\text{C}_6\text{H}_5\text{XCF}_3 \text{ or } \text{C}_6\text{H}_5\text{XCH}_3) - q(\text{C}_6\text{H}_6)$ . The quantum-chemical calculation for benzene at MP2/6-31G(d) level followed by NBO analysis of the obtained wavefunction shows that the natural charges comprise –0.2350 electron on each carbon and 0.2350 electron on each hydrogen atom of the molecule. The calculated values of  $\Delta q$  are presented in Table 5. Increase in the negative value (or decrease in the positive value) of  $\Delta q$  corresponds to rise in the atomic charge compared to the corresponding atom (C or H) in unsubstituted benzene. Due to the higher electronegativity of oxygen compared to carbon, the C<sub>6</sub>H<sub>5</sub>OCH<sub>3</sub> and C<sub>6</sub>H<sub>5</sub>OCF<sub>3</sub> molecules bear excess electronic density on the oxygen and low electronic density on the C<sub>i</sub> atom of the aromatic ring. The values of  $\Delta q_{C_i}$  (~0.5 electron) found for C<sub>6</sub>H<sub>5</sub>OCH<sub>3</sub> and C<sub>6</sub>H<sub>5</sub>OCF<sub>3</sub> indicates the strong acceptor

Table 4  
Atomic natural charges ( $q$ , electron) for C<sub>6</sub>H<sub>5</sub>XY (X = O or S; Y = CH<sub>3</sub> or CF<sub>3</sub>)

$\varphi^\circ$	$q_X$	$q_{C_y}$	$\sum(qY)/3$	$q_{C_i}$	$q_{C_o}$	$q_{H_o}$	$q_{C_m}$	$q_{H_m}$	$q_{C_p}$	$q_{H_p}$
C <sub>6</sub> H <sub>5</sub> OCF <sub>3</sub>										
0	–0.5710	1.3638	–0.3615	0.2877	–0.2896, –0.2598	0.2563, 0.2549	–0.2193, –0.2233	0.2425, 0.2424	–0.2412	0.2411
45	–0.5774	1.3629	–0.3625	0.2798	–0.2797, –0.2577	0.2588, 0.2550	–0.2195, –0.2222	0.2428, 0.2428	–0.2380	0.2415
90	–0.5826	1.3613	–0.3636	0.2660	–0.2583, –0.2575	0.2561, 0.2558	–0.2216, –0.2216	0.2431, 0.2431	–0.2349	0.2416
C <sub>6</sub> H <sub>5</sub> SCF <sub>3</sub>										
0	0.2903	0.9402	–0.3588	0.2057	–0.2482, –0.2429	0.2489, 0.2437	–0.2175, –0.2209	0.2420, 0.2421	–0.2360	0.2405
45	0.2736	0.9415	–0.3604	–0.2166	–0.2401, –0.2365	0.2525, 0.2481	–0.2169, –0.2195	0.2427, 0.2426	–0.2311	0.2409
90	0.2583	0.9433	–0.3620	–0.2302	–0.2270, –0.2270	0.2523, 0.2522	–0.2186, –0.2186	0.2433, 0.2433	–0.2267	0.2414
C <sub>6</sub> H <sub>5</sub> OCH <sub>3</sub>										
0	–0.5306	–0.3095	0.2092	0.3061	–0.3134, –0.2625	0.2385, 0.2470	–0.2270, –0.2325	0.2354, 0.2368	–0.2515	0.2358
45	–0.5489	–0.3037	0.2065	0.2990	–0.2908, –0.2628	0.2386, 0.2473	–0.2295, –0.2315	0.2361, 0.2370	–0.2466	0.2362
90	–0.5656	–0.2965	0.2037	0.2896	–0.2677, –0.2673	0.2448, 0.2450	–0.2321, –0.2321	0.2372, 0.2372	–0.2403	0.2368
C <sub>6</sub> H <sub>5</sub> SCH <sub>3</sub>										
0	0.3038	–0.8395	0.2483	–0.2060	–0.2586, –0.2380	0.2340, 0.2413	–0.2245, –0.2273	0.2369, 0.2383	–0.2420	0.2368
45	0.2742	–0.8266	0.2447	–0.2109	–0.2426, –0.2363	0.2367, 0.2455	–0.2255, –0.2259	0.2367, 0.2388	–0.2367	0.2374
90	0.2507	–0.8183	0.2421	–0.2171	–0.2324, –0.2324	0.2459, 0.2461	–0.2270, –0.2270	0.2391, 0.2391	–0.2310	0.2379

Table 5

The differences of atomic charges in C<sub>6</sub>H<sub>5</sub>XCF<sub>3</sub>, C<sub>6</sub>H<sub>5</sub>XCH<sub>3</sub> (X = O and S) and in unsubstituted benzene [ $\Delta q = q(\text{C}_6\text{H}_5\text{XCF}_3 \text{ or } \text{C}_6\text{H}_5\text{XCH}_3) - q(\text{C}_6\text{H}_6)$ ]

$\varphi^\circ$	$\Delta q C_i$	$\Delta q C_o$	$\Delta q H_o$	$\Delta q C_m$	$\Delta q H_m$	$\Delta q C_p$	$\Delta q H_p$
C <sub>6</sub> H <sub>5</sub> OCF <sub>3</sub>							
0	0.5227	-0.0546, -0.0248	0.0213, 0.0199	0.0157, 0.0117	0.0075, 0.0074	-0.0062	0.0061
45	0.5148	-0.0447, -0.0227	0.0238, 0.0200	0.0155, 0.0128	0.0078, 0.0078	-0.0030	0.0065
90	0.5010	-0.0233, -0.0225	0.0211, 0.0208	0.0134, 0.0134	0.0081, 0.0081	0.0001	0.0066
C <sub>6</sub> H <sub>5</sub> SCF <sub>3</sub>							
0	0.0293	-0.0132, -0.0079	0.0139, 0.0087	0.0175, 0.0141	0.0070, 0.0071	-0.0010	0.0055
45	0.0184	-0.0051, -0.0015	0.0051, 0.0015	0.0181, 0.0155	0.0077, 0.0076	0.0039	0.0059
90	0.0048	0.0080, 0.0080	0.0173, 0.0172	0.0164, 0.0164	0.0083, 0.0083	0.0083	0.0064
C <sub>6</sub> H <sub>5</sub> OCH <sub>3</sub>							
0	0.5411	-0.0784, -0.0276	0.0035, 0.0120	0.0080, 0.0025	0.0004, 0.0018	-0.0165	0.0008
45	0.5340	-0.0558, -0.0278	0.0036, 0.0123	0.0055, 0.0035	0.0011, 0.0020	-0.0116	0.0012
90	0.5246	-0.0327, -0.0323	0.0098, 0.0100	0.0029, 0.0029	0.0022, 0.0022	-0.0053	0.0018
C <sub>6</sub> H <sub>5</sub> SCH <sub>3</sub>							
0	0.0290	-0.0236, -0.0030	-0.0010, 0.0063	0.0105, 0.0077	0.0019, 0.0033	-0.0070	0.0018
45	0.0241	-0.0076, -0.0013	0.0017, 0.0105	0.0095, 0.0091	0.0017, 0.0038	-0.0017	0.0024
90	0.0179	0.0026, 0.0026	0.0109, 0.0111	0.0080, 0.0080	0.0041, 0.0041	0.0040	0.0029

effect of the OCH<sub>3</sub> and OCF<sub>3</sub> groups on the C<sub>i</sub> atoms. The C<sub>6</sub>H<sub>5</sub>SCH<sub>3</sub> and C<sub>6</sub>H<sub>5</sub>SCF<sub>3</sub> molecules have lower electron density on the C<sub>i</sub> atom. In the C<sub>6</sub>H<sub>5</sub>OCH<sub>3</sub> molecule in the planar conformation the most favourable conditions for interaction of the LP2 on the oxygen atom with  $\pi^*$ -orbitals of the ring are realized. For this reason, the values of  $\Delta q$  on the atoms C<sub>o</sub> (-0.0784 and -0.0275 electron) and C<sub>p</sub> (-0.0165 electron) are more negative than in other compounds. If the equilibrium in C<sub>6</sub>H<sub>5</sub>OCF<sub>3</sub> is shifted to near-orthogonal conformations ( $\varphi = 45-90^\circ$ ), the excess of electron density is 0.02 electron on the CO and 0.003 electron on the C<sub>p</sub> atoms (compared to unsubstituted benzene). In the orthogonal conformation, the charge on the C<sub>p</sub> atom is practically equal to the charge on carbons in unsubstituted benzene. The SCH<sub>3</sub> group in C<sub>6</sub>H<sub>5</sub>SCH<sub>3</sub> has a donor effect on the *ortho*- and *para*-position of the ring only in near-planar conformations or when free rotation is present ( $\varphi = 45^\circ$ ). In the orthogonal form SCH<sub>3</sub>-groups acts on C<sub>o</sub> and C<sub>p</sub> as an acceptor. In orthogonal C<sub>6</sub>H<sub>5</sub>SCF<sub>3</sub>, the SCF<sub>3</sub> substituent also displays an acceptor effect on the C<sub>o</sub> and C<sub>p</sub> atoms. In all molecules considered, the XCF<sub>3</sub> and XCH<sub>3</sub> groups produce the acceptor effect on the C<sub>s</sub> and on hydrogen atoms of the aromatic ring.

On nitration, C<sub>6</sub>H<sub>5</sub>OCF<sub>3</sub> gives a *para*-isomer, whereas C<sub>6</sub>H<sub>5</sub>SCF<sub>3</sub> is nitrated both at the *ortho*- and *para*-position [12,13]. As is evident from the charge distribution, the charge on the C<sub>o</sub> and C<sub>p</sub> atoms in C<sub>6</sub>H<sub>5</sub>OCF<sub>3</sub> is higher than in C<sub>6</sub>H<sub>5</sub>SCF<sub>3</sub>, that is, the SCF<sub>3</sub> group deactivates the benzene ring towards electrophilic agents to a greater extent than does the OCF<sub>3</sub> group. The selective *para*-directive effect of the OCF<sub>3</sub> group is evidently associated with a small rotational barrier in the C<sub>6</sub>H<sub>5</sub>OCF<sub>3</sub> molecule. As OCF<sub>3</sub> freely rotates about the C<sub>sp<sup>2</sup></sub>-O bond, it sterically hampers attack on the *ortho*-position of the aromatic ring and the agent is directed to the *para*-position (though the electronic charge on the C<sub>o</sub> atoms is higher than on the C<sub>p</sub>).

The C<sub>6</sub>H<sub>5</sub>SCF<sub>3</sub> molecule has a stable orthogonal conformation with almost equal charges on the C<sub>o</sub> and C<sub>p</sub> atoms and for this reason one should expect an equimolar ratio of the *ortho*- and *para*-isomers. However, despite the orthogonal conformation of the molecule, there is a steric shielding of the *ortho*-position with the bulky CF<sub>3</sub> group and, as a result, C<sub>6</sub>H<sub>5</sub>SCF<sub>3</sub> gives 65% of the *para*- and 35% of the *ortho*-isomer.

Thus, the distribution of electron density in aromatic ring of the C<sub>6</sub>H<sub>5</sub>XCF<sub>3</sub> and C<sub>6</sub>H<sub>5</sub>XCH<sub>3</sub> molecules is governed by two major factors, the electronic nature of the substituents XCF<sub>3</sub> or XCH<sub>3</sub> and the molecular conformation and agree well with reactivity.

## 7. NMR spectra

Table 6 lists chemical shifts of <sup>1</sup>H, <sup>13</sup>C and <sup>19</sup>F signals in NMR spectra of compounds C<sub>6</sub>H<sub>5</sub>XCF<sub>3</sub> and C<sub>6</sub>H<sub>5</sub>XCH<sub>3</sub>. It is known that in <sup>13</sup>C NMR spectra of monosubstituted benzenes, the change in shielding of the carbon nuclei in the *para*-position is correlated with relative charge on this atom [49–51]. To derive the relationship of the  $\delta C_p$  and the natural charge on the *para*-carbon atoms we calculated, in the MP2/6-31G(d) approximation, the values of  $q C_p$  in molecules C<sub>6</sub>H<sub>5</sub>R with the following R (the charge in *e* on C<sub>p</sub> is given in parentheses): NH<sub>2</sub> (-0.2620), F (-0.2466), Cl (-0.2384), Br (-0.2369), CF<sub>3</sub> (-0.2265), and CN (-0.2261). The values of  $q C_p$  for the above compounds, unsubstituted benzene and the C<sub>6</sub>H<sub>5</sub>OCH<sub>3</sub> and C<sub>6</sub>H<sub>5</sub>SCF<sub>3</sub> molecules in the preferred conformations correlate reasonably well with  $\delta C_p$  for the compounds [24,51] though the linear equation:

$$\delta C_p = 231.29(\pm 3.55) + 436.43(\pm 14.83)q C_p,$$

$$r = 0.996, \quad S_0 = 0.5216, \quad n = 9$$



Table 6

Chemical shifts of  $^1\text{H}$  ( $\delta\text{H}$ , ppm from TMS),  $^{13}\text{C}$  ( $\delta\text{C}$ , ppm from TMS) and  $^{19}\text{F}$  ( $\delta\text{F}$ , ppm from  $\text{CFCl}_3$ ) signals in NMR spectra in  $\text{C}_6\text{H}_5\text{XCF}_3$  and  $\text{C}_6\text{H}_5\text{XCH}_3$  ( $\text{X} = \text{O}$  or  $\text{S}$ )<sup>a,b</sup>

Compound	$\delta\text{C}_i$	$\delta\text{C}_o$	$\delta\text{C}_m$	$\delta\text{C}_p$	$\delta\text{C}_\gamma$	$\delta\text{F}$	$\delta\text{H}_\gamma$
$\text{C}_6\text{H}_5\text{OCH}_3$	150.42 (21.92)	122.17 (−6.33)	131.31 (2.81)	128.50 (0.00)	121.81	−57.54	
$\text{C}_6\text{H}_5\text{SCH}_3$	159.93 (31.43)	114.01 (−14.49)	129.47 (0.97)	120.61 (−7.89)	54.59		3.76 <sup>c</sup>
$\text{C}_6\text{H}_5\text{OCF}_3$	124.71 (−3.79)	137.07 (8.57)	130.51 (2.01)	131.93 (3.43)	130.75	−43.18	
$\text{C}_6\text{H}_5\text{SCF}_3$	139.35 (10.85)	126.90 (−1.60)	129.45 (0.95)	125.43 (−3.07)	15.40		2.44 <sup>c</sup>

<sup>a</sup> In acetone- $\text{d}_6$  solution.

<sup>b</sup> The differences of chemical shifts on the atoms in the molecules and in unsubstituted benzene are given in parentheses.

<sup>c</sup> Taken from [54].

The correlation plot is shown in Fig. 4. If the values of  $\delta\text{C}_p$  for the conformationally non-rigid molecules  $\text{C}_6\text{H}_5\text{SCH}_3$  and  $\text{C}_6\text{H}_5\text{OCF}_3$  are placed on the curve and the points are projected onto the coordinate axis, it is possible to estimate the charge on the  $\text{C}_p$  atoms in the molecules. The results indirectly suggest that the conformational equilibrium in

solution is shifted to orthogonal for  $\text{C}_6\text{H}_5\text{OCF}_3$  or to planar for  $\text{C}_6\text{H}_5\text{SCH}_3$ . Qualitatively similar, but less pronounced dependence of  $\delta\text{C}$  and  $q\text{C}$  is also observed, for other atoms of the aromatic ring. However, the magnetic shielding of nuclei is determined by many factors, of which the charge distribution is not always the determining one. The  $^{13}\text{C}$

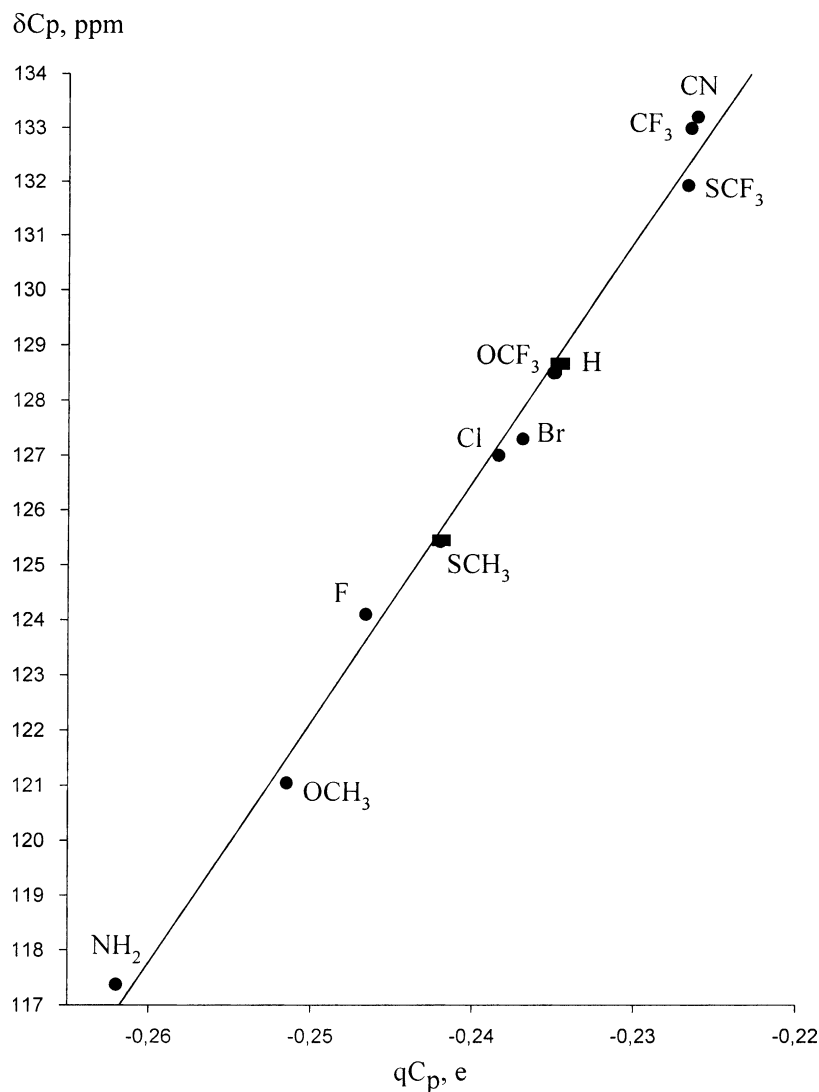


Fig. 4. Correlation between NMR  $^{13}\text{C}$  ( $\delta\text{C}_p$ , ppm) chemical shifts of *para*-carbon atoms and natural charges ( $q\text{C}_p$ , electron) on these atoms in monosubstituted benzenes.

nuclei in the CH<sub>3</sub> group is substantially more shielded than in the CF<sub>3</sub> group which corresponds to the charges on these atoms, but, with a sulfur atom substituted for oxygen, the shielding of <sup>13</sup>C nuclei in the CF<sub>3</sub> group decreases, where in the CH<sub>3</sub> group it rises. In the CH<sub>3</sub> group, the change in δC parallels the change in the charge. In the CF<sub>3</sub> group, the shielding of the <sup>13</sup>C nucleus and the charge on this atom change in opposite directions as O is substituted by S. Such a substitution also causes opposite trends in the changes of shielding of the <sup>19</sup>F and <sup>1</sup>H nuclei in the CF<sub>3</sub> and CH<sub>3</sub> groups. In this case the changes in shielding of protons are opposite to the changes in the charge. The values of δO (ppm relative to H<sub>2</sub>O; given in parentheses is the width at half-height) were found to be 112 ppm ( $\nu_{1/2} = 160$  Hz) for C<sub>6</sub>H<sub>5</sub>OCF<sub>3</sub> and 49 ppm ( $\nu_{1/2} = 190$  Hz) for C<sub>6</sub>H<sub>5</sub>OCH<sub>3</sub>. They are in good agreement with the reported values of 113 ppm [25] (C<sub>6</sub>H<sub>5</sub>OCF<sub>3</sub>), 45 ppm [52] and 48 ppm [53] (C<sub>6</sub>H<sub>5</sub>OCH<sub>3</sub>). The spin–spin coupling constants (<sup>1</sup>J<sub>CF</sub>) are negative [49]. In absolute value the <sup>1</sup>J<sub>CF</sub> in C<sub>6</sub>H<sub>5</sub>OCF<sub>3</sub> (–255.8 Hz) is smaller than in C<sub>6</sub>H<sub>5</sub>SCF<sub>3</sub> (–309.5 Hz). The <sup>1</sup>J<sub>CH</sub> in C<sub>6</sub>H<sub>5</sub>OCH<sub>3</sub> (143.5 Hz [54]) is larger than in C<sub>6</sub>H<sub>5</sub>SCH<sub>3</sub> (140 Hz [54]) and the range of its variation in the thio analogs is substantially less wide.

## 8. Conclusion

Investigation of the potential function of internal rotation about the C<sub>sp<sup>2</sup></sub>–X bonds in the C<sub>6</sub>H<sub>5</sub>XCF<sub>3</sub> and C<sub>6</sub>H<sub>5</sub>XCH<sub>3</sub> (X = O, S) molecules at the MP2/6-31G(d) level of theory gives new information about spatial structures of these compounds. The C<sub>6</sub>H<sub>5</sub>OCH<sub>3</sub> molecule is planar while C<sub>6</sub>H<sub>5</sub>SCF<sub>3</sub> is orthogonal. C<sub>6</sub>H<sub>5</sub>SCH<sub>3</sub> and C<sub>6</sub>H<sub>5</sub>OCF<sub>3</sub> have either free rotation of the fragments about the C<sub>sp<sup>2</sup></sub>–O bond or molecular fragment motion with a large amplitude in the broad flat bottomed potential well and a small rotational barrier. Owing to the small rotation barrier, experimental conditions (solvent, temperature) can affect on the profile of the internal rotation potential function and markedly shift the conformation equilibrium in the C<sub>6</sub>H<sub>5</sub>OCF<sub>3</sub> and C<sub>6</sub>H<sub>5</sub>SCH<sub>3</sub> molecules. The hybridization of two LP's on the atoms O and S is different. One of the LP's comprises a hybrid orbital (with ≈40% of s-component for O and ≈70% for S). The second LP is a practically pure p-orbital. In near-planar conformations the second LP of O and S atoms interact strongly with π\*-system of the ring. In similar conformations such interaction is somewhat more effective for XCH<sub>3</sub> than for XCF<sub>3</sub> provided that the heteroatoms are the same. Interaction is substantially lowered in non-planar conformations and practically vanishes in the orthogonal form. However, in non-planar conformations, the hybridized LP of the oxygen or sulfur atom is involved in the interaction with π\*-orbitals, which peaks in the orthogonal form. For oxygen such interaction is more effective than for sulfur. The electron density in the aromatic ring is determined not only by differences in electronic effects of the XCF<sub>3</sub>

and XCH<sub>3</sub> groups but also and to a large measure by their conformation.

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